

Supporting information for

Critical Role of Surface Energy in Guiding Crystallization of Solution Coated Conjugated Polymer Thin Films

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Supplementary Note: Substrate Surface Energy Determination

Based on the Young's Equation¹, the contact angle of a pure liquid droplet on a smooth and chemically homogenous solid surface at thermodynamic equilibrium is defined as the balance of three interfacial tensions (Figure S1 and Equation S1):

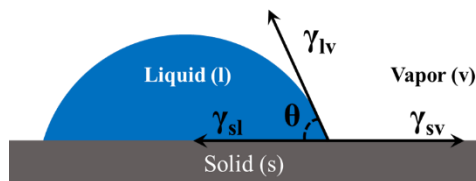


Figure S1. Contact angle at the three-phase contact line.

$$\gamma_{sv} = \gamma_{sl} + \gamma_{lv} \cos \theta \quad (\text{S1})$$

where γ_{lv} , γ_{sl} and γ_{sv} are liquid-vapor, solid-liquid and solid-vapor interfacial free energies respectively, and θ is the static contact angle (often quasi-static advancing contact angle is used²). Liquid-air interfacial free energy of the pure probing liquid (γ_{lv}) and contact angle (θ) can be measured using the pendant and sessile drop methods and axisymmetric drop shape analysis using standard Goniometers.³ However, one equation is insufficient for determining two unknowns, γ_{sv} and γ_{sl} . It has been shown experimentally that γ_{sv} , γ_{sl} and γ_{lv} are not independent states of the system, but are related via an equation of state⁴⁻⁶. In other words, γ_{sl} can be fully determined by γ_{lv} and γ_{sv} (assuming the system is in thermodynamic equilibrium)^{2, 4}:

$$\gamma_{sl} = f(\gamma_{sv}, \gamma_{lv}) \quad (\text{S2})$$

Specifically, we employ the following equation-of-state relation^{2, 7} in our work:

$$\gamma_{sl} = \gamma_{lv} + \gamma_{sv} - 2\sqrt{\gamma_{lv}\gamma_{sv}}(1 - \beta(\gamma_{lv} - \gamma_{sv})^2) \quad (\text{S3})$$

where β is an empirical constant. The $(\gamma_{lv} + \gamma_{sv} - 2\sqrt{\gamma_{lv}\gamma_{sv}})$ term comes from the geometric mean combining rule, which is rooted in London theory of dispersion and assumes one dominant type of interactions between the liquid and the solid molecules, or like pairs⁸. However, the combining rule overestimates the interaction between two dissimilar pairs (unlike pairs) and an empirical parameter $(1-\beta(\gamma_{lv}-\gamma_{sv})^2)$ is added to correct for this deviation from experimental results, while maintaining the symmetry of combining rules.² By combining Equation S3 with the Young's equation (Equation S1), $\cos \theta$ can be expressed as,

$$\cos \theta = -1 + 2 \sqrt{\frac{\gamma_{sv}}{\gamma_{lv}}} (1 - \beta(\gamma_{lv} - \gamma_{sv})^2) \quad (\text{S4})$$

By measuring contact angles (θ) with a series of probing liquids with known but diverse surface energies (γ_{lv}), substrate surface energy (γ_{sv}) and β specific for each substrate can be calculated using the least squares regression method.⁹ Table S1 summarizes eleven probing liquids selected in this work (chemically inert with respect to substrates within the time frame of contact angle measurements), containing a diverse range of chemistries, intermolecular forces and surface energies (measured in this work or taken from external references).

Table S1. Probing liquids and their surface tensions at ambient temperature.

Probing Liquid	Surface Tension (mN/m)	Reference
Dodecane	25.4±0.2	¹⁰
1-Pentanol	26.0	¹¹
Hexadecane	27.6	¹¹
1-Ethyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide ([EMIM][TFSI])	36.9 ±0.7	This work
Dimethyl sulfoxide	44.0±1.5	This work

		and ¹¹
Ethylene glycol	47.7±0.7	This work and ¹¹
Diiodomethane	50.0	¹¹
Formamide	59.1	¹¹
Glycerol	65.0	¹¹
Water	72.7	¹¹
Mercury	470	¹¹

The measurement of θ in the Young's equation is crucial; in accurate θ measurements or the choice of the wrong type of θ (advancing vs. receding, static vs. dynamic) can lead to erroneous results. It has been shown that the quasi-static advancing contact angle, θ_a , is the best approximation to the thermodynamic equilibrium static contact angle¹², which was carefully measured in this study by slowly supplying excess liquid to the sessile drop (the three-phase contact line will move outward during this process). The final volume of the drop was set to be less than 3 μL to avoid the deformation of drop shape by gravitational force. A Rame-Hart Model 250 Standard Goniometer equipped with drop shape analysis software was employed in this work for measuring the contact angles. For each probing liquid and substrate pair, at least 10 different advancing contact angles were collected on new surfaces making sure that the substrates are chemically homogenous and our results are reproducible. The average advancing contact angles thus measured are summarized in Table S2.

Table S2. Advancing contact angles of the probing liquids on various substrates.

	OTS	FPTS	PTS	PVPHDA	Bare	Plasma
Dodecane	32.9±0.3	-	-	-	-	-
1-Pentanol	37.1±0.2	-	-	-	-	-
Hexadecane	41.1±0.2	-	-	-	-	-
[EMIM][TFSI]	75.1±0.5	31.0±0.5	-	-	-	-
Dimethyl	74.3±0.8	43.2±2.0	29.7±0.9	-	-	-

sulfoxide						
Ethylene glycol	85.4±0.4	62.3±0.3	54.1±0.6	-	-	-
Diiodomethane	76.2±0.7	61.0±0.3	46.7±0.7	35.6±1.7	-	-
Formamide	86.4±1.8	65.9±0.1	59.9±0.9	-	27.1±0.2	-
Glycerol	97.1±0.6	82.0±0.4	74.1±0.6	53.7±2.5	43.4±2.2	12.4±2.3
Water	112.3±0.3	91.1±0.7	80.2±1.5	63.2±3.8	52.0±0.7	9.9±1.7
Mercury	-	-	-	-	-	128.0±1.1

Note that γ_{lv} of the probing liquids usually need to be higher than the solid surface energy γ_{sv} to avoid total wetting conditions ($\theta=0$), under which the Young's equation is no longer valid. Strictly, the total wetting condition is when $\gamma_{lv} + \gamma_{sl} < \gamma_{sv}$, determined by the spreading coefficient. But when $\gamma_{lv} < \gamma_{sv}$, the contact angle can be too small to measure accurately.

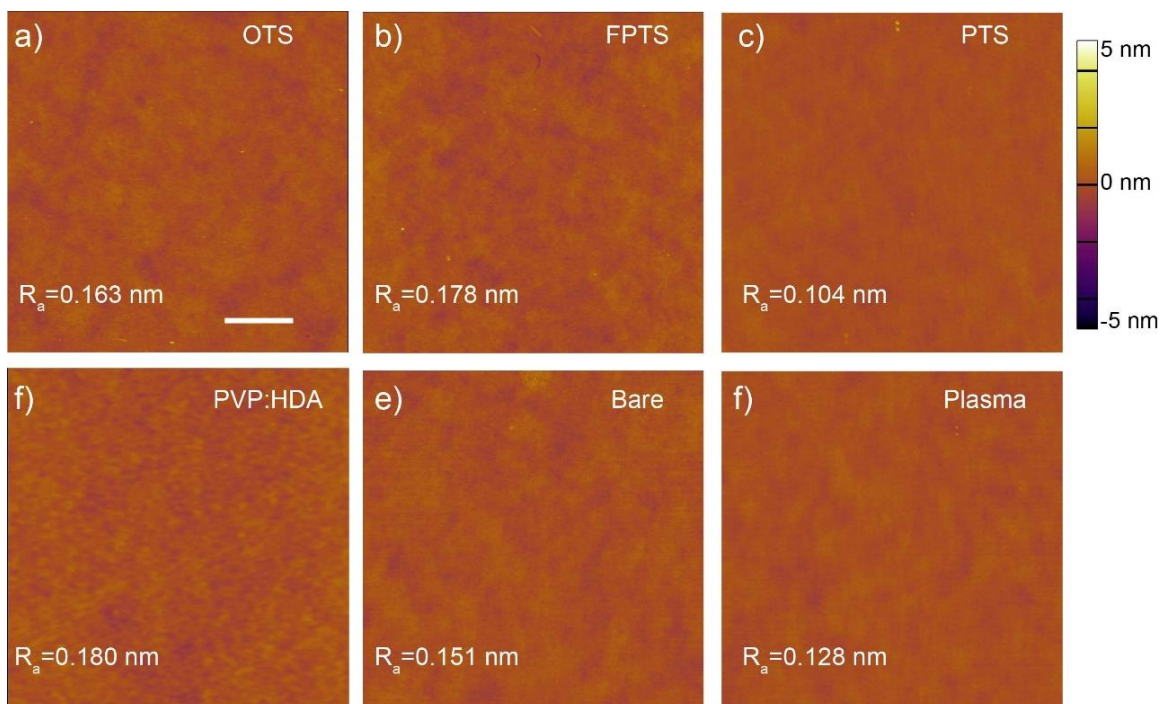


Figure S2. AFM images of the top-surface morphology for (a) OTS, (b) FPTs, (c) PTS, (d) PVP:HDA, (e) bare and (d) plasma treated SiO₂ surface. The scale bare is 1 μ m. The mean roughness (R_a) is labeled on the lower left of each image.

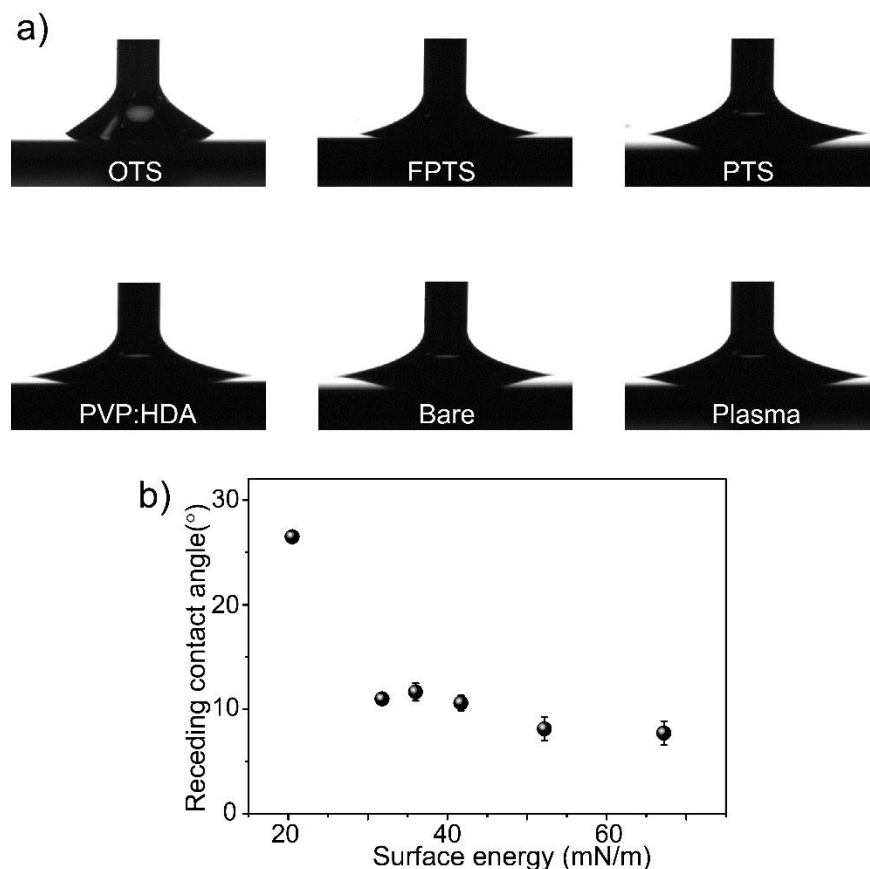


Figure S3. Receding contact angles of chloroform on various substrates. (a) Optical images of receding meniscus of chloroform on various substrates. (b) Receding contact angle as a function of substrate surface energy. The receding contact angles of chloroform on different substrates, which approximate the dynamic contact angle during meniscus-guided coating, were measured during solvent evaporation at room temperature. The receding contact angles for all surfaces fall within a narrow range of 15° and only OTS substrate has a relatively higher value possibly due to the crystalline nature of the OTS SAM as well as the non-polarity imparted by the alkyl chains. Therefore, the meniscus shape for DPP2T-TT/chloroform solution and evaporation rates are comparable for most of the substrates except for OTS. In cases with comparable evaporation rates and wetting properties, the OSC film morphology differed drastically. For instance, even

though the receding contact angle is similar for chloroform on PTS and PVP:HDA, the degree of crystallinity dramatically decreased comparing films on PTS vs. PVP:HDA. This result validates that the surface energy primarily impacts polymer crystallization, not the ink wetting and drying behaviors. In addition, solvent evaporation rate (Q_{evap}) in meniscus-guided coating configuration can be described from the evaporation regime model by Le Berre et. al. derived from mass balance equation¹³: $Q_{\text{evap}} = h \rho L v / C$; where h is the film thickness, ρ is the solute density, L is the meniscus width, v is the coating speed and C is the solution concentration. We controlled polymer film thickness to be ~ 30 nm across all substrates by slightly changing concentration within the narrow range of 3.8 ± 0.5 mg/mL for accurate comparison of UV-vis, C-POM and GIXD data. Therefore, Q_{evap} should be comparable across substrates based on the equation above.

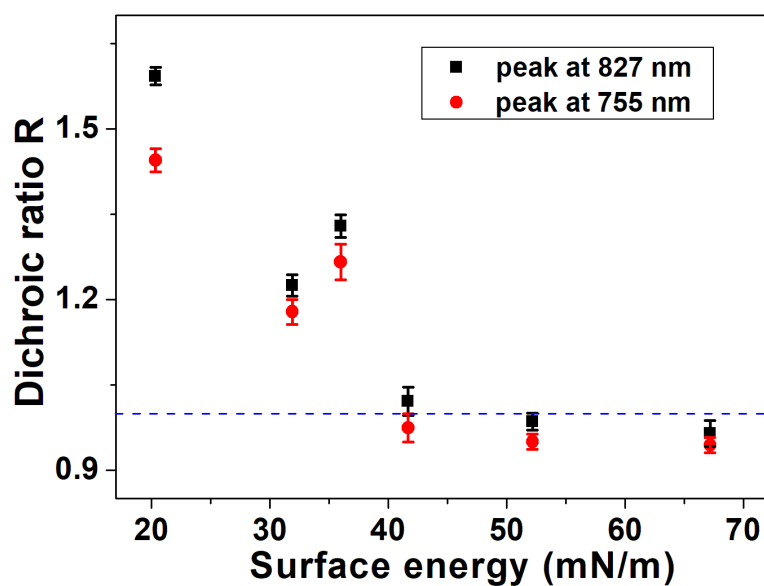


Figure S4. UV-vis dichroic ratio R ($R = I_{\text{perp}}/I_{\text{par}}$) of DPP2T-TT films as a function of substrate surface energy. R was calculated from the peak absorbance at 755 nm and 827 nm. The dash line indicates $R=1$.

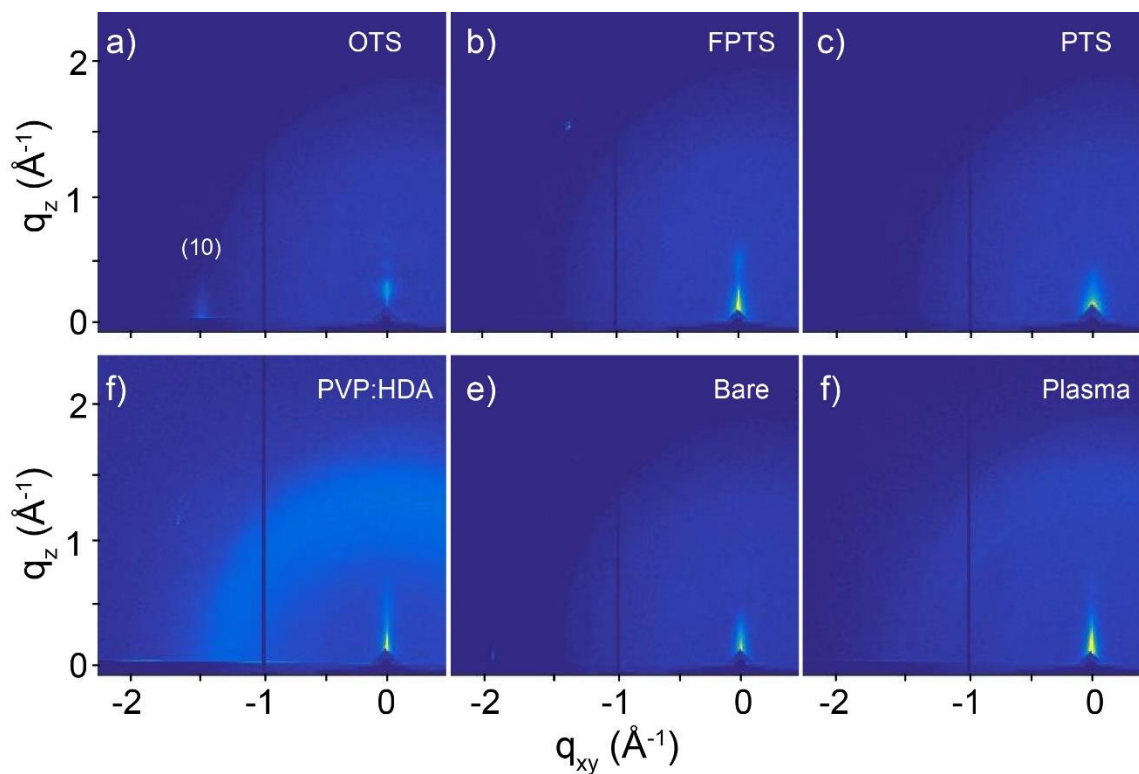


Figure S5. GIXD images of (a) OTS, (b) FPTS, (c) PTS, (d) PVP:HDA, (e) bare and (f) plasma treated Si substrate before coating the polymer semiconductor film. The diffraction of substrates demonstrated crystalline structures of the OTS monolayer, while others shown the amorphous packing. The OTS-SAM layer showed a peak at $q_{xy} = 1.49 \text{ \AA}^{-1}$, corresponding to the (10) diffraction of the hexagonal crystalline lattice.¹⁴⁻¹⁵

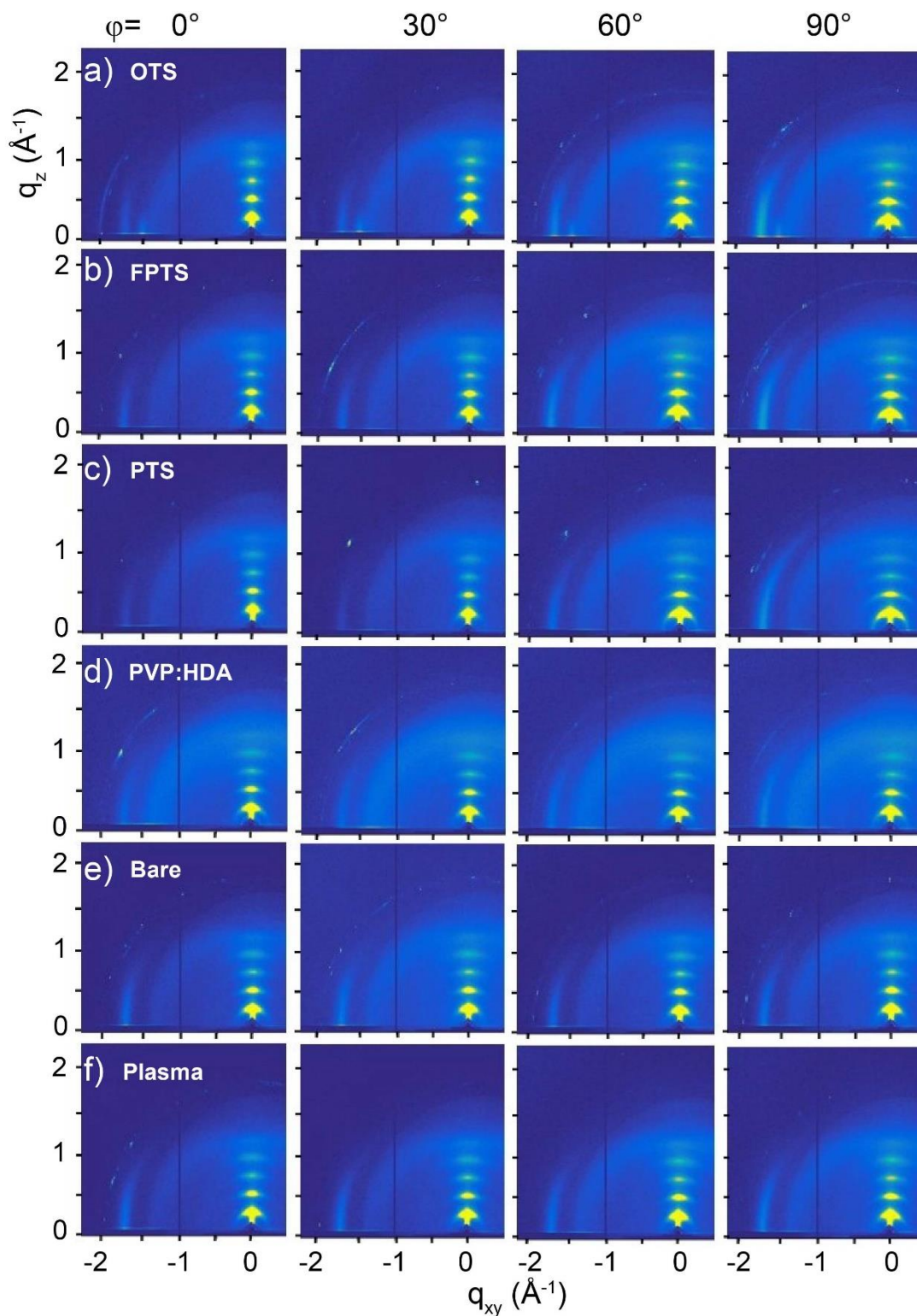


Figure S6. GIXD images of DPP2T-TT films on (a) OTS, (b) FPTs, (c) PTS, (c) PVP:HDA, (e) bare and (f) plasma treated Si substrate, at $\varphi = 0^\circ, 30^\circ, 60^\circ$ and 90° . The

angle φ is the in-plane rotation angle of the sample, and 0° is defined as when the coating direction is aligned with the incident beam. Note that geometrical correction was not applied to these raw images to better visualize the out-of-plane peaks. We did apply geometrical correction during analysis of the diffraction patterns.

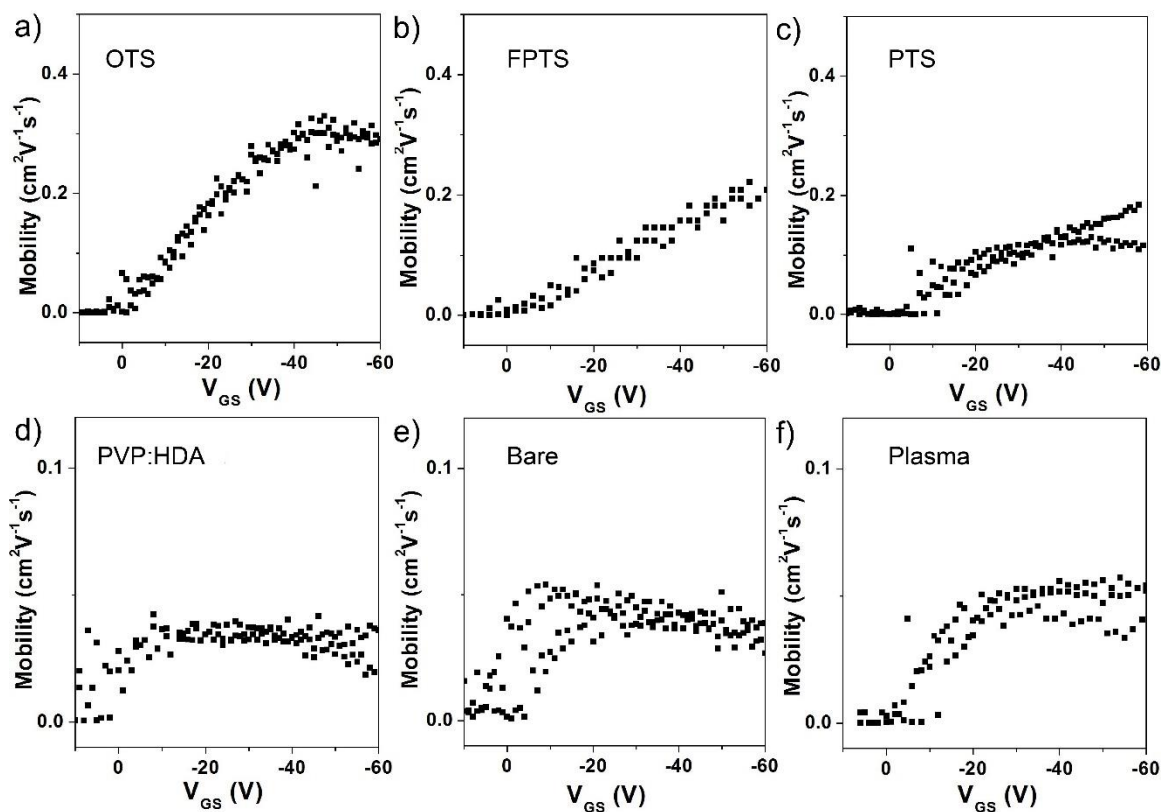


Figure S7. V_{GS} dependent apparent mobilities of DPP2T-TT OFETs with TGTC configuration in the saturation regime. The semiconducting layers were deposited on (a) OTS, (b) FPTS, (c) PTS, (d) PVP:HDA, (e) bare and (f) plasma treated SiO_2 substrates. $V_{DS} = -60$ V; the mobility is calculated per 2V for forward and reverse scanning.

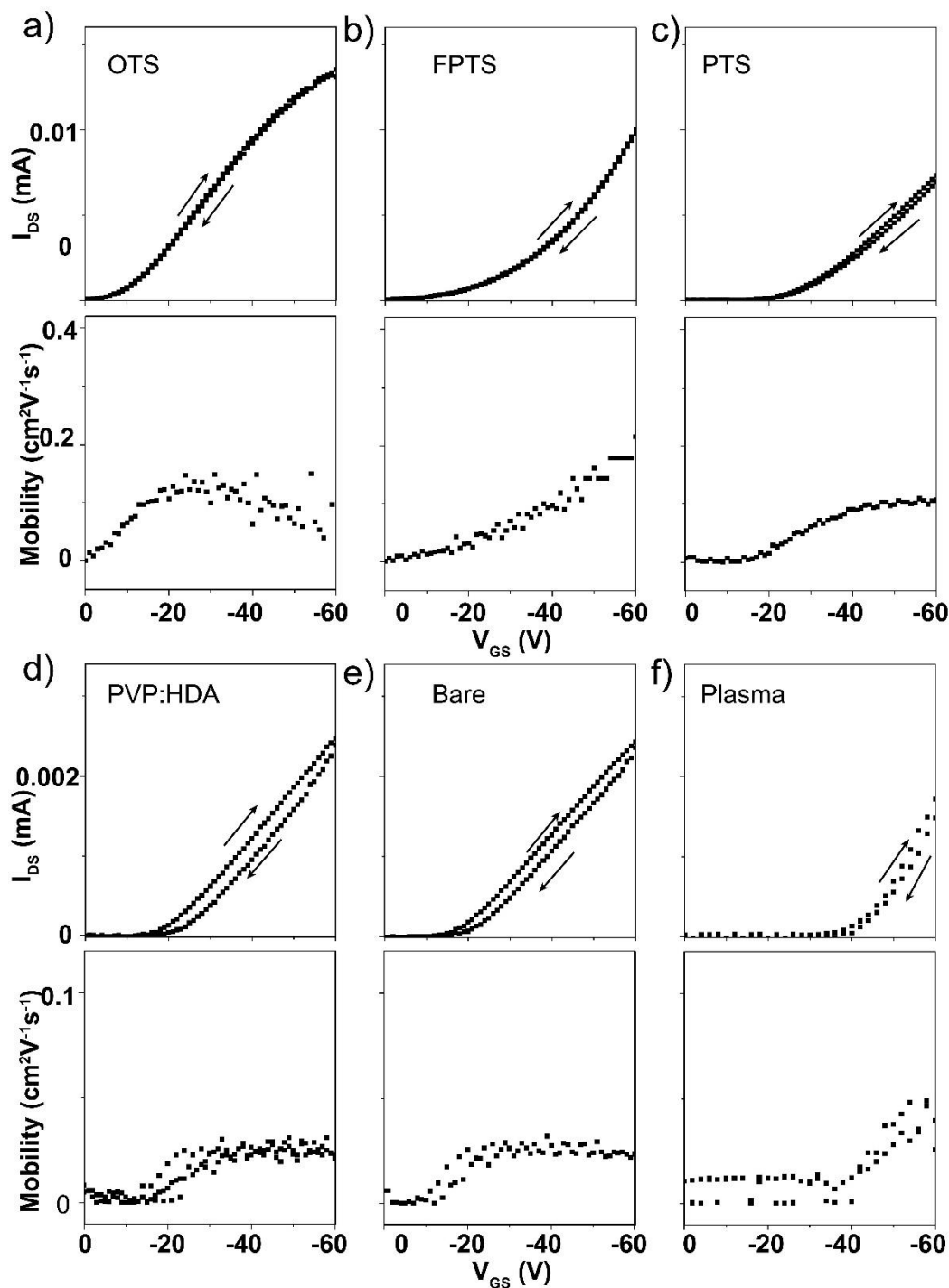


Figure S8. Transfer curves (Top) and V_{GS} dependent apparent mobilities (bottom) of DPP2T-TT TGTC OFETs in the linear regime, obtained at $V_{DS} = -10$ V. The semiconducting layers were deposited on (a) OTS, (b) FPTS, (c) PTS, (d) PVP:HDA, (e) bare and (f) plasma treated SiO_2 substrates.

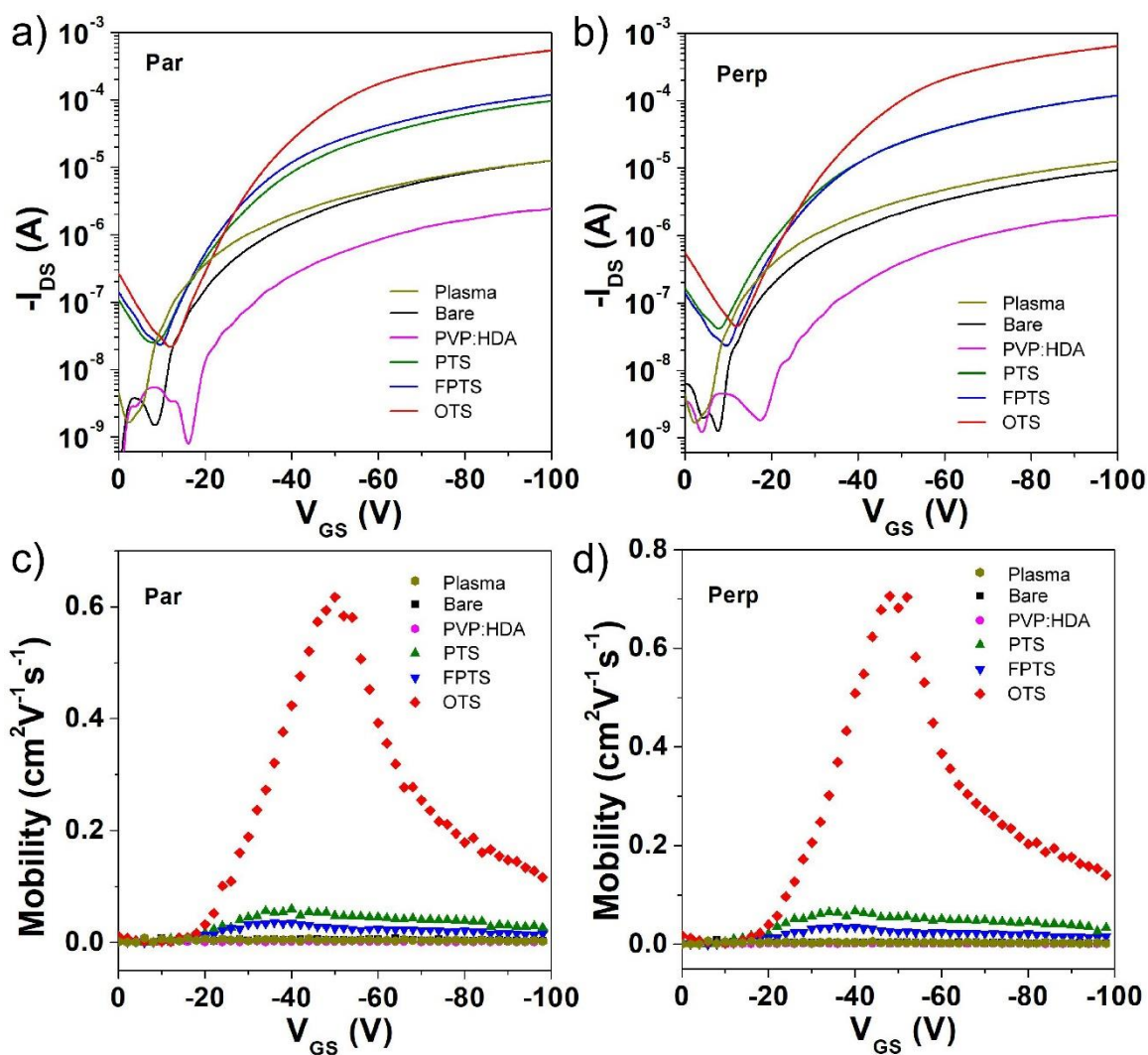


Figure S9. (a, b) Transfer curves and (c, d) mobility as a function of gate voltage for DPP2T-TT OFETs with BGTC configuration, in the saturation regime with $V_{DS} = -100$ V for the measurements. ‘Par’ and ‘Perp’ denote the case when the channel is parallel and perpendicular to the coating direction, respectively.

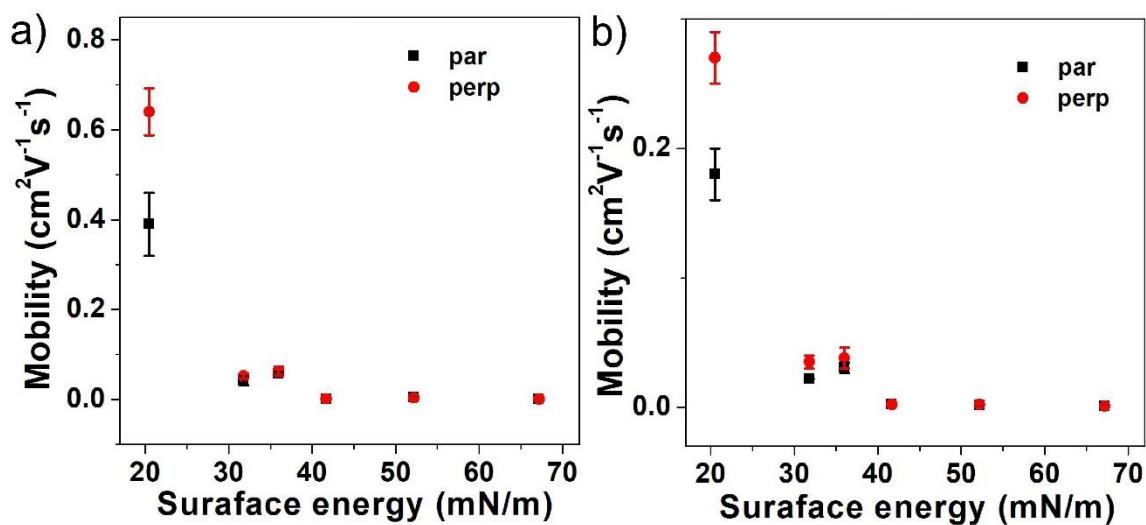


Figure S10. Apparent mobilities of BGTC OFETs as a function of surface energy. The mobility was calculated in the saturation regime in the (a) low ($0 < |V_{GS}| < 60$ V) and (b) high gate voltage ($60 < |V_{GS}| < 100$ V) regions.

Table S3. Summary of device characteristics for DPP2T-TT OFETs with TGTC and BGTC structure based on DPP2T-TT films on various substrate.

Device structure	Substrate	Surface energy (mN/m)	Channel length direction to coating	Mobility ($\text{cm}^2\text{V}^{-1}\text{s}^{-1}$)		V_{th} (V)	Log (I_{on}/I_{off})
				Low V_{GS}	High V_{GS}		
TGTC	OTS	20.5±1.4	Par	0.22±0.06		-14.5±5.4	4.7±0.9
			Perp	0.29±0.08		-15.1±3.5	4.8±0.2
	FPTS	31.8±1.5	Par	0.09±0.04		-10.6±3.2	4.1±0.3
			Perp	0.18±0.04		-12.5±1.5	4.2±0.2
	PTS	36.0±1.6	Par	0.08±0.03		-5.2±2.4	4.1±0.1
			Perp	0.14±0.02		-7.5±3.1	4.2±0.5
	PVP:HDA	41.7±1.2	Par	0.03±0.001		9.0±1.0	4.2±0.5
			Perp	0.04±0.004		-8.3±3.8	3.9±0.6
	Bare	52.2±1.0	Par	0.05±0.010		-5.4±4.5	3.8±0.6
			Perp	0.04±0.002		-7.3±5.9	5.2±0.1
	Plasma	67.2±2.4	Par	0.04±0.006		-9.0±1.1	4.2±0.2
			Perp	0.05±0.003		-12.0±0.8	4.3±0.1
BGTC	OTS	20.5±1.4	Par	0.49±0.07	0.16±0.02	-18.3±2.0	3.5±0.6
			Perp	0.74±0.02	0.18±0.01	-20.1±0.2	4.3±0.2
	FPTS	31.8±1.5	Par	0.04±0.003	0.02±0.003	-15.0±3.1	4.0±0.8
			Perp	0.05±0.001	0.03±0.005	-14.5±1.6	3.6±0.3
	PTS	36.0±1.6	Par	0.06±0.004	0.03±0.004	-21.0±2.5	3.8±0.8
			Perp	0.07±0.008	0.04±0.008	-18.0±1.0	3.6±0.3
	PVP:HDA	41.7±1.2	Par	0.001±0.0002	0.002±0.0002	-13.0±2.0	2.9±0.5
			Perp	0.001±0.0003	0.002±0.0003	-17.5±2.5	3.2±0.3
	Bare	52.2±1.0	Par	0.005±0.0002	0.004±0.0002	-13.5±1.5	3.1±0.4
			Perp	0.004±0.0001	0.006±0.0001	-9.0±1.0	3.2±0.2
	Plasma	67.2±2.4	Par	0.005±0.0002	0.004±0.0002	-15.5±2.5	3.5±0.4
			Perp	0.004±0.0001	0.005±0.0003	-11.0±3.7	3.1±0.5

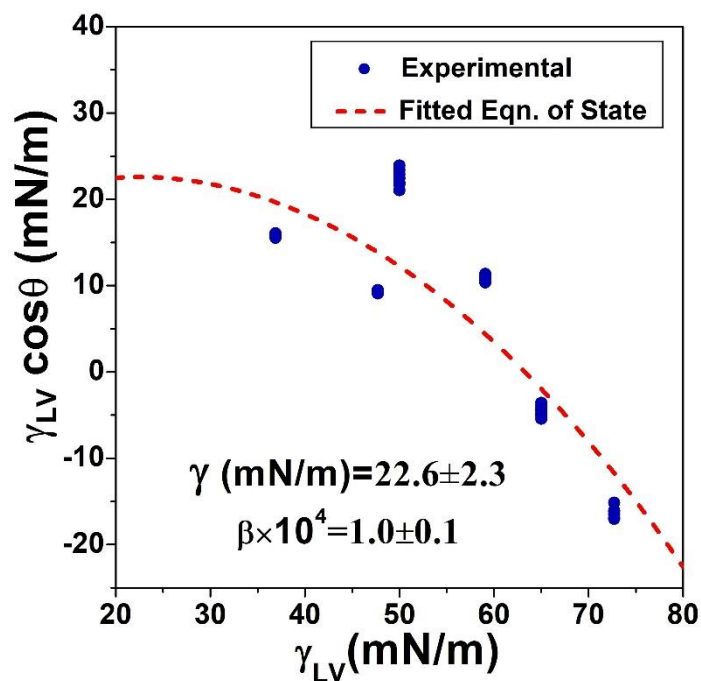


Figure S11. Determination of DPP2T-TT surface energy via the equation of state approach. The plot shown is the $\gamma_{LV} \times \cos \theta$ vs. γ_{LV} plot comparing experimental data and best-fitted curve obtained from the equation of state.

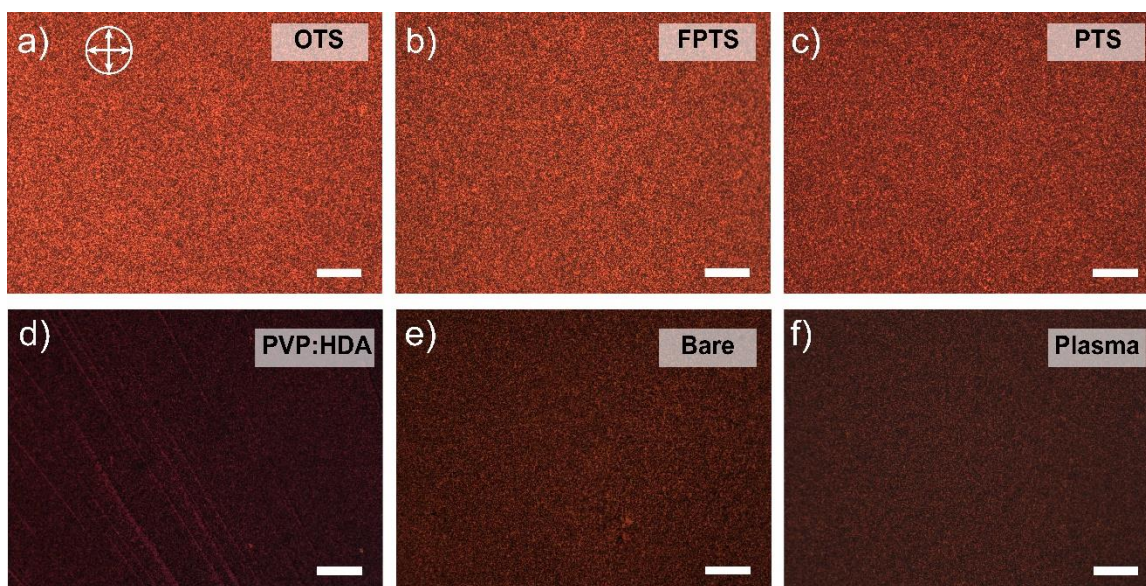


Figure S12. C-POM images for spin-coated DPP2T-TT thin films on (a) OTS, (b) FPTS, (c) PTS, (d) PVP:HDA, (e) bare and (f) plasma treated substrate, taken at the same exposure time. All scale bars are 100 μm .

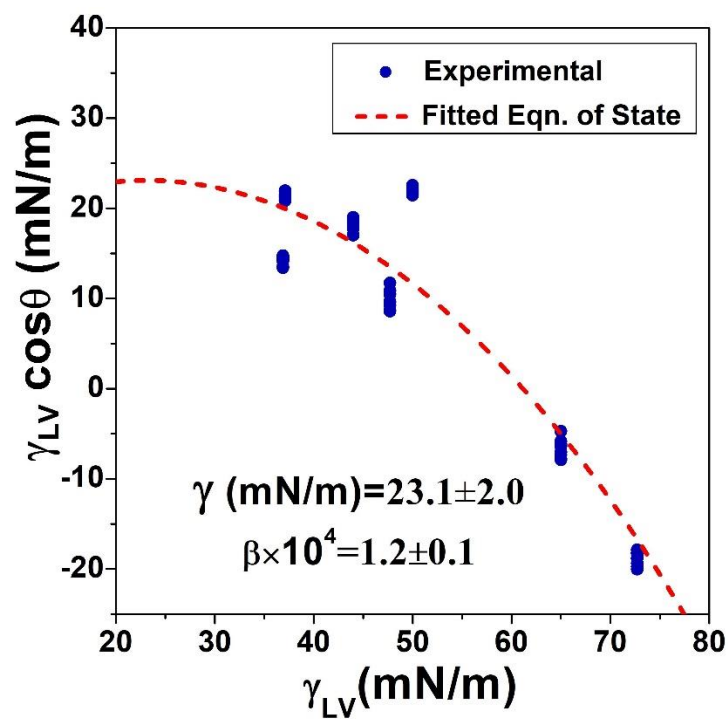


Figure S13. Determination of PII-2T surface energy via the equation of state approach. The plot shown is the $\gamma_{LV} \times \cos \theta$ vs. γ_{LV} plot comparing experimental data and best-fitted curve obtained from the equation of state.

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